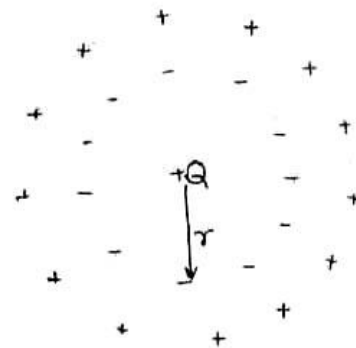


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**Debye Length**

Consider a charge +Q placed in plasma. It will attract electrons all around and the ions will be repelled by this. In equilibrium, the charge produces electric field around it then

$\vec{E} = -\nabla \phi$  (i.e. gradient of scalar potential  $\phi$ )  
 ..... (i)



At a distance r from test charge Q there is a charge particle, then potential energy

P.E = -e  $\phi$  for electrons

P.E = e  $\phi$  for ions

Boltzmann's law says that electrons will have a tendency to go in a region of small potential energy, then

Electron density  $n_e = n_0 e^{\frac{e\phi}{KT_e}}$

Ion density  $n_i = n_0 e^{\frac{-e\phi}{KT_i}}$

Where  $n_0$  = equilibrium density of plasma electrons which is same as plasma ions

For simplicity  $T_e = T_i = T$

Poisson's equation in electrostatics

$\epsilon_0 \nabla \cdot \vec{E} = \rho$

Where  $\rho$  = charge density

$\rho = n_i e - n_e e$

$\rho = (n_i - n_e)e$

$\therefore \nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$

$\nabla \cdot \vec{E} = \frac{e}{\epsilon_0} (n_i - n_e)$

From eqn. (i)  $\nabla \cdot (-\nabla\phi) = \frac{e}{\epsilon_0} (n_i - n_e)$

$$\nabla^2\phi = \frac{e}{\epsilon_0} (n_e - n_i)$$

If  $\frac{e\phi}{KT} \ll 1$ , then

$$n_e \simeq n_0 \left(1 + \frac{e\phi}{KT}\right)$$

$$n_i \simeq n_0 \left(1 - \frac{e\phi}{KT}\right)$$

$$\therefore \nabla^2\phi = \frac{en_0}{\epsilon_0} \left[1 + \frac{e\phi}{KT} - 1 + \frac{e\phi}{KT}\right]$$

$$\nabla^2\phi = \frac{2n_0e^2}{\epsilon_0KT} \phi$$

The potential due to a charge  $q$  as a function of distance. So in that case

$$\frac{1}{r^2} \frac{\delta}{\delta r} \left( r^2 \frac{\delta\phi}{\delta r} \right) = \frac{2n_0e^2}{\epsilon_0KT} \phi$$

The coefficient in the right hand side that multiplies by the quantity  $\phi$  has one upon length square then

$$\lambda_D^2 = \frac{\epsilon_0KT}{2n_0e^2}$$

$\lambda_D$  has a dimension of length. It depends upon temperature and electron density.

$$\text{Poisson's equation can be written as } \frac{1}{r^2} \frac{\delta}{\delta r} \left( r^2 \frac{\delta\phi}{\delta r} \right) = \frac{\phi}{\lambda_D^2}$$

Now solve this equation

By the fact that the potential of charge placed in free space goes as

$$\phi = \frac{1}{r} F(r) \text{ Where } F(r) \text{ is some constant}$$

Substitute in above equation

$$\frac{d^2F}{dr^2} = \frac{F}{\lambda_D^2}$$

The solution of above equation is simple exponential solution

$$F = C_1 e^{\frac{-r}{\lambda_D}}$$

It has a physical argument that the potential must decrease as more away from charge.

Now put F in  $\phi$

$$\phi = \frac{C_1}{r} e^{-\frac{r}{\lambda_D}}$$

To evaluate  $C_1$  the screening effect of electrons or ions will not be effective because in an electron cloud or ion cloud, the electric field is zero

If  $r \ll \lambda_D$ , then potential

$$\phi = \frac{Q}{4\pi\epsilon_0 r}, \quad C_1 = \frac{Q}{4\pi\epsilon_0}$$

Potential can be written as

$$\phi = \frac{Q}{4\pi\epsilon_0 r} e^{-\frac{r}{\lambda_D}}$$

The term  $e^{-\frac{r}{\lambda_D}}$  is an extra screening factor and the potential will fall off quite rapidly when  $r > \lambda_D$

So  $\lambda_D =$  Debye length

$$\lambda_D = \sqrt{\frac{\epsilon_0 kT}{2n_0 e^2}}$$

The significance is that the electrostatic potential of a charge is felt over a distance of the order of  $\lambda_D$  beyond which it is screened out.

So the dimension of plasma should be bigger than Debye length.