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Langevin's theory of Paramagnetism -

Consider a molecular gas containing N molecules or atoms per unit volume at a temperature T . Each atom or molecule is assumed to possess a permanent magnetic moment \mathbf{m} . If all the dipoles are along the applied magnetic field, the magnetization \mathbf{M} will be $N\mathbf{m}$. Actually the dipoles are pointing in all directions. The magnetic potential energy of the dipole, which is inclined at an angle θ with the applied field \mathbf{B} , is given by

$$U = -\mathbf{m} \cdot \mathbf{B} = -mB \cos \theta. \quad \dots(1)$$

According to the Boltzmann distribution law, the relative probability of finding a dipole making an angle θ with the \mathbf{B} - direction is $e^{-U/kT}$, where k is the Boltzmann constant. Thus the number of dipoles dN , out of total number N , inclined between angles θ and $\theta + d\theta$ with the magnetic field \mathbf{B} will be given by

$$\begin{aligned} dN &\propto e^{-U/kT} d\Omega \\ &= C e^{-U/kT} d\Omega, \end{aligned} \quad \dots(2)$$

where C is a constant of proportionality and $d\Omega$ is the solid angle subtended by the cones of semi-vertical angles θ and $\theta + d\theta$ and given by

$$d\Omega = 2\pi \sin \theta d\theta$$

$$\begin{aligned} \therefore dN &= C e^{-U/kT} \cdot 2\pi \sin \theta d\theta = C \cdot 2\pi e^{mB \cos \theta / kT} \sin \theta d\theta. \\ &= 2\pi C e^{\alpha \cos \theta} \sin \theta d\theta, \end{aligned} \quad \dots(3)$$

where

$$\alpha = mB/kT.$$

Integration of Eq. (103) will give the total number of dipoles per unit volume N . Thus

$$N = \int_0^\pi 2\pi C e^{\alpha \cos \theta} \sin \theta d\theta = \frac{4\pi C \sinh \alpha}{\alpha}$$

or

$$C = N\alpha/4\pi \sinh \alpha. \quad \dots(4)$$

\therefore

$$dN = \frac{N\alpha}{2 \sinh \alpha} e^{\alpha \cos \theta} \sin \theta d\theta. \quad \dots(5)$$

The component of each dipole moment parallel to \mathbf{B} is $m \cos \theta$, hence the total magnetic moment per unit volume

$$\begin{aligned}
 M &= \int_0^\pi m \cos \theta dN = \frac{mN\alpha}{2 \sinh \alpha} \int_0^\pi e^{\alpha \cos \theta} \sin \theta \cos \theta d\theta \\
 &= \frac{mN\alpha}{2 \sinh \alpha} \left[\frac{e^{\alpha \cos \theta}}{\alpha^2} - \frac{\cos \theta e^{\alpha \cos \theta}}{\alpha} \right]_0^\pi \\
 &= Nm (\coth \alpha - 1/\alpha). \quad \dots(6)
 \end{aligned}$$

Since Nm is the saturation value of the magnetic moment, say M_0 , hence

$$M = M_0 (\coth \alpha - 1/\alpha) = M_0 L(\alpha). \quad \dots(7)$$

This bracket term $(\coth \alpha - 1/\alpha)$ is known as *Langevin function*, $L(\alpha)$. It is plotted as a function of α , in Fig. 13.25. This curve indicates that for large values of α the function tends to unity, i.e., $M \rightarrow M_0$. It is the case when all the atomic dipoles are parallel to \mathbf{B} . For small values of α , the curve is linear, as for smaller values of α the Langevin function

$$L(\alpha) = \coth \alpha - 1/\alpha = \alpha/3 = mB/3kT.$$

For most cases of paramagnetics, $\alpha \ll 1$ at ordinary temperature even for \mathbf{B} as large as 1 tesla, therefore \mathbf{M} can be approximated by the lower values of α , as

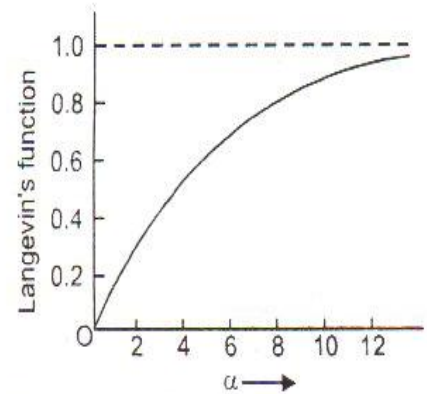
$$\mathbf{M} = \frac{M_0 m \mathbf{B}}{3kT} = \frac{Nm^2 \mathbf{B}}{3kT}. \quad \dots(8)$$

For a paramagnetic substances, \mathbf{M} is very small positive value, hence $\mathbf{B} = \mu_0 (\mathbf{H} + \mathbf{M}) \approx \mu_0 \mathbf{H}$. Therefore, we have

$$\mathbf{M} = Nm^2 \mu_0 \mathbf{H} / 3kT. \quad \dots(9)$$

For small values of α and thus for higher values of T

$$\begin{aligned}
 \text{Susceptibility } \chi &= \frac{M}{H} = \frac{M_0 (\coth \alpha - 1/\alpha)}{H} = \frac{m^2 N \mu_0 H}{3kTH} \\
 &= \frac{\mu_0 M_0^2}{3NkT} = \frac{\text{const}}{T}. \quad \dots(10)
 \end{aligned}$$



Variation of Langevin function.

Here the constant is called the *Curie constant* and this relation is known as *Curie law*. Curie law is well verified experimentally and is very reasonable as at higher temperatures the thermal agitation opposes the dipole alignments along the \mathbf{B} -direction. It breaks down at low temperatures. The temperature below which this law ceases can reach as low as 1.3°K , for some paramagnetic substances. At a very low temperature $\alpha \gg 1$ and $M = Nm$, which is independent of magnetic field \mathbf{B} . Thus M takes its saturation value M_0 .