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Subject — Physics.

Title of the topic — Motion in a central field.

Name of Program — B.Sc. (H) PHYSICS

Part-I, Paper-I (Group-B)

Two body problem Reduced to one body problem

Central force — Central force b/w two particles is one which is directed along the line joining the two particles and whose magnitude is fn. of distance  $r$  b/w them.

i.e.  $\vec{F} = F(r) \hat{r}$ .

Conclusion:

①  $\vec{L} = 0$

②  $\vec{J} = \text{const.}$

③  $\vec{J}$  is  $\perp$  to  $\vec{r}$  &  $\vec{p}$   $\therefore \vec{J} = \vec{r} \times \vec{p}$ .

$\Rightarrow \vec{r} \cdot \vec{J} = \vec{r} \cdot (\vec{r} \times \vec{p}) = 0$

$\Rightarrow$  motion of particle in a plane

General form of Central force.

or

Inverse  $n^{\text{th}}$  power law:

$$\vec{F} = \frac{C}{r^n} \hat{r} \quad \forall C \text{ is a const.}$$

$$\therefore -\frac{dU}{dr} = \frac{C}{r^n}$$

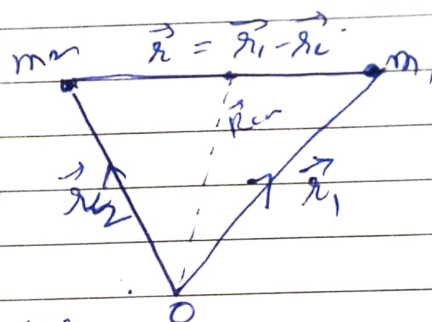
$U \rightarrow$  P.E. of the particle

$$\Rightarrow U = - \int \frac{C}{r^n} dr$$

Consider the motion of a system of 2 particles subjected to a central force. it is a 2 body problem.

we want to reduce it to a 1 body problem

let us consider two particles of masses  $m_1$  &  $m_2$  of position vectors  $\vec{r}_1$  &  $\vec{r}_2$  w.r.t O.



$\therefore$  Central force exerted by the 2nd particle on the 1st,

$$\vec{F}_{12} = m_1 \left( \frac{d^2 \vec{r}_1}{dt^2} \right) \Rightarrow \frac{d^2 \vec{r}_1}{dt^2} = \frac{\vec{F}_{12}}{m_1}$$

Central force exerted by 1st particle on the 2nd,

$$\vec{F}_{21} = m_2 \left( \frac{d^2 \vec{r}_2}{dt^2} \right) \Rightarrow \frac{d^2 \vec{r}_2}{dt^2} = \frac{\vec{F}_{21}}{m_2}$$

Since  $\vec{F}_{12} = -\vec{F}_{21} = \vec{F} = F \hat{r} \rightarrow$  Let the central force is gravitational force.

$$\therefore \frac{d^2 \vec{r}_1}{dt^2} = \frac{\vec{F}}{m_1} \quad \& \quad \frac{d^2 \vec{r}_2}{dt^2} = \frac{-\vec{F}}{m_2}$$

$$\therefore \frac{d^2 \vec{r}_1}{dt^2} - \frac{d^2 \vec{r}_2}{dt^2} = \frac{\vec{F}}{m_1} + \frac{\vec{F}}{m_2} = \left( \frac{1}{m_1} + \frac{1}{m_2} \right) \vec{F}$$

$$\frac{d^2 \vec{r}_1}{dt^2} - \frac{d^2 \vec{r}_2}{dt^2} = \left( \frac{1}{m_1} + \frac{1}{m_2} \right) (-G \frac{m_1 m_2}{r^2} \hat{r})$$

$$\Rightarrow \frac{d^2 \vec{r}}{dt^2} = - \frac{1}{r^2} \left( \frac{m_1 m_2}{\frac{1}{m_1} + \frac{1}{m_2}} \right) \hat{r}$$

$$\Rightarrow \mu \frac{d^2 \vec{r}}{dt^2} = - \frac{G m_1 m_2}{r^2} \hat{r}$$

Reduced mass of the system

$$\left\{ \begin{aligned} \therefore \frac{1}{\mu} &= \frac{1}{m_1} + \frac{1}{m_2} \\ \therefore \mu &= \frac{m_1 m_2}{m_1 + m_2} \end{aligned} \right.$$

Since we get an eqn. involving only one position vector  $\vec{r}$   $\therefore$  it is one body problem.



8

3

$$\Rightarrow \frac{d^2 \vec{r}}{dt^2} = -G(m_1 + m_2) \frac{\vec{r}}{r^3}$$

Tuesday

or

$\Rightarrow$  It is the eqn. of motion of particle of unit mass at a vector distance  $\vec{r}$  from a fixed mass  $(m_1 + m_2)$  ~~located at focus of a ellipse~~

Week 28  
18/11/162003  
1 yr

Note:  $\vec{r}_1 - \vec{r}_2 = \vec{r}$  &  $m_1 \vec{r}_1 + m_2 \vec{r}_2 = (m_1 + m_2) \vec{R}_{cm}$

$$\therefore \vec{r}_1 = \vec{R}_{cm} + \frac{m_2 \vec{r}}{m_1 + m_2} = \vec{R}_{cm} + \frac{\mu \vec{r}}{m_1}$$

$$\vec{r}_2 = \vec{R}_{cm} - \frac{m_1 \vec{r}}{m_1 + m_2} = \vec{R}_{cm} - \frac{m_1}{m_1 + m_2} \vec{r} = \vec{R}_{cm} - \frac{\mu}{m_2} \vec{r}$$

Case: ① if  $m_1 = m_2 = m \therefore \mu = \frac{m}{2}$

② if  $m_1 \ll m_2 \therefore \mu = \frac{m_1 m_2}{m_2 \left( \frac{m_1}{m_2} + 1 \right)}$

$$\approx m_1 \left( 1 + \frac{m_1}{m_2} \right)^{-1}$$

$$\approx m_1 \left( 1 - \frac{m_1}{m_2} \right)$$

Reduced mass of H-atom ( $1e^- + 1p$ )

$$m_p \approx 1836 m_e$$

$$\mu_H = \frac{m_p m_e}{m_p + m_e}$$

$$m_e \ll m_p \therefore \mu_H \approx m_e \left( 1 - \frac{m_e}{m_p} \right)$$

$$\approx m_e \left( 1 - \frac{1}{1836} \right) \approx m_e$$

$\therefore$  reduced mass is ~~value~~ smaller mass

Reduced mass of positronium atom ( $1 \text{ positron} + 1e^-$ )

$$\therefore \mu_p = \frac{m_e m_e}{m_e + m_e} = \frac{1}{2} m_e$$

( $e^+$  &  $e^-$ )

Note  $\therefore$  reduced mass of positronium atom

$$= \frac{1}{2} m_e$$

" " Hydrogen atom

UL

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4

So we may determine the motion of  $m_1$  relative to  $m_2$  by taking  $m_2$  to lie at fixed origin of inertial frame & replace  $m_1$  by  $\mu$ .

(41)

Wednesday

9

Week 28  
190-175

9.00

Angular mom. for equivalent 1 body problem

10.00

$$\vec{J} = \vec{r} \times \vec{p}$$

11.00

$$= \vec{r} \times \mu \frac{d\vec{r}}{dt} = \vec{r} \times \mu \left[ \dot{r} \hat{r} + r \dot{\theta} \hat{\theta} \right]$$

12.00

$$= r \hat{r} \times \mu \left[ \dot{r} \hat{r} + r \dot{\theta} \hat{\theta} \right]$$

$\left\{ \begin{array}{l} r, \theta, z \end{array} \right.$  cylindrical coords

$$= 0 + \mu r^2 \dot{\theta} \hat{z}$$

1.00

$$\boxed{\vec{J} = \mu r^2 \omega \hat{z}}$$

2.00

in 2 body problem

3.00

$$\vec{J} = \vec{r}_1 \times m_1 \frac{d\vec{r}_1}{dt} + \vec{r}_2 \times m_2 \frac{d\vec{r}_2}{dt}$$

4.00

for 1D

Since  $\vec{J} = \vec{r} \times \mu \frac{d\vec{r}}{dt} = (\vec{r}_1 - \vec{r}_2) \times \left( \frac{m_1 m_2}{m_1 + m_2} \right) \left[ \frac{d\vec{r}_1}{dt} - \frac{d\vec{r}_2}{dt} \right]$

Went to C.M. frame of ref.

6.00

$$\therefore \vec{r} = \vec{r}_1 - \vec{r}_2$$

Since  $\vec{R}_{cm} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2}$

7.00

$$\Rightarrow (m_1 + m_2) \dot{\vec{R}}_{cm} = m_1 \dot{\vec{r}}_1 + m_2 \dot{\vec{r}}_2$$

$$\Rightarrow m_1 \dot{\vec{r}}_1 = -m_2 \dot{\vec{r}}_2 \quad \left\{ \because \dot{\vec{R}}_{cm} = 0 \right.$$

8.00

$$\Rightarrow \frac{d\vec{r}_1}{dt} = -\frac{m_2}{m_1} \frac{d\vec{r}_2}{dt} \quad \text{--- (2)}$$

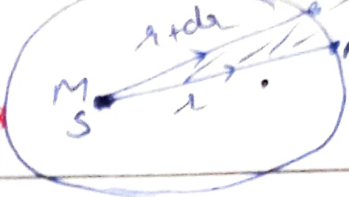
9.00

$$\text{f } \frac{d\vec{r}_2}{dt} = -\frac{m_1}{m_2} \frac{d\vec{r}_1}{dt} \quad \text{--- (3)}$$

$\therefore$  Put (2) & (3) in (1)

$$\vec{J} = (\vec{r}_1 - \vec{r}_2) \left( \frac{m_1 m_2}{m_1 + m_2} \right) \left[ \frac{m_1}{m_2} \frac{d\vec{r}_1}{dt} - \frac{m_2}{m_1} \frac{d\vec{r}_2}{dt} \right]$$





For the planetary motion

from Δ SAB

$$\frac{\text{area swept}}{\text{Time}} = \frac{\frac{1}{2} \vec{r} \times d\vec{r}}{dt}$$

Since  $I = mrv$   
 $\Rightarrow \frac{I}{m} = rv$

$$\begin{aligned} \frac{d\vec{s}}{dt} &= \frac{1}{2} rv \\ &= \frac{1}{2} \frac{I}{m} \\ &= \text{constant} \\ &= r^2 \dot{\theta} \end{aligned}$$

area vel.

In Polar Coordinates

$$x = r \cos \theta$$

$$y = r \sin \theta$$

$$\vec{V} = v_1 \hat{r} + v_2 \hat{\theta}$$

radial comp

In x dir

$$v_1 \cos \theta - v_2 \sin(90 - \theta) = \frac{dx}{dt}$$

$$v_1 \cos \theta - v_2 \sin \theta = \frac{d}{dt} (r \cos \theta) = \dot{r} \cos \theta - r \sin \theta \dot{\theta}$$

$$\therefore v_1 = \dot{r} \quad \& \quad v_2 = r\dot{\theta}$$

In y dir

$$v_1 \sin \theta - v_2 \sin(90 - \theta) = \frac{dy}{dt} = \frac{d}{dt} (r \sin \theta)$$

$$\Rightarrow v_1 \sin \theta + v_2 \cos \theta = \dot{r} \sin \theta + r \cos \theta \dot{\theta}$$

$$\Rightarrow v_1 = \dot{r} \quad \& \quad v_2 = r\dot{\theta}$$

now  $\vec{a} = a_1 \hat{r} + a_2 \hat{\theta}$

In x dir

$$a_1 \cos \theta - a_2 \sin(90 - \theta) = \frac{dv_x}{dt}$$

$$= \frac{d}{dt} [\dot{r} \cos \theta - r \sin \theta \dot{\theta}]$$

$$\therefore a_1 \cos \theta - a_2 \sin \theta = \dot{x} \cos \theta - \dot{x} \sin \theta - \dot{x} \sin \theta \dot{\theta} - \dot{x} \cos \theta (\dot{\theta})^2 - \dot{x} \sin \theta \ddot{\theta}$$

10.00

radial comp

$$\therefore a_r = \ddot{r} - r(\dot{\theta})^2$$

$$f \quad a_2 = 2 \ddot{x} + \dot{x}^2 = \ddot{x} + 2 \dot{x} \dot{x}$$

12.00 Tangential comp.  $= \frac{1}{r} \frac{d}{dt} (r^2 \dot{\theta}) = 0 \quad \left\{ \because r^2 \dot{\theta} = \text{const} \right.$   
 $\Rightarrow \hat{a}_\theta = \frac{1}{r} \frac{d}{dt} (r^2 \dot{\theta}) \hat{\theta} = (\ddot{\theta} - \frac{2}{r} \dot{r} \dot{\theta}) \hat{\theta}$

12.00  $\therefore \vec{a} = a_1 \hat{x} + a_2 \hat{y} = (\lambda' - \lambda_0'^2) \hat{x}$

1.00  $\therefore$  In planetary motion no tangential comp of  $acc^n$  occurs  
ie total  $acc^n$  is ~~only~~ along radial dir<sup>n</sup> and no comp  
2.00 of force will occur along tang. dir<sup>n</sup>.

of force will occurs along tangential

Consider the motion of particle of mass  $m$  around  $\odot$  such that distance b/w them is  $r$ .

$$\therefore \vec{F} = - \frac{G M m}{r^2} \hat{r} = m \vec{a} = 4 \frac{d^2 r}{dt^2}$$

$$4.00 \Rightarrow -\frac{GMm}{r^2} \hat{r} = m \left[ \frac{d^2 \mathbf{r}}{dt^2} - r \left( \frac{d\phi}{dt} \right)^2 \hat{r} \right]$$

$$\Rightarrow \frac{d^2 r}{dt^2} - r \left( \frac{d\theta}{dt} \right)^2 = -\frac{GM}{r^2}$$

$$\Rightarrow r^3 \frac{d^2 r}{dt^2} - r^4 \left( \frac{d\theta}{dt} \right)^2 = -GM r \quad \left\{ \begin{array}{l} \text{multiply} \\ r^3 \text{ on both} \\ \text{sides} \end{array} \right.$$

$$\Rightarrow r^3 \frac{d^2 r}{dr^2} - r^4 \left( \frac{d\theta}{dt} \right)^2 =$$

$$\Rightarrow r^3 \frac{d^2 r}{dt^2} - h^2 = -GMm \quad \text{--- (1)} \quad \left[ \begin{array}{l} \text{Let } h = r^2 \dot{\theta} = \text{const.} \\ \& u = \frac{1}{r} \end{array} \right]$$

13 Sunday *Now*  $h = r^2 \dot{\phi} \Rightarrow \dot{\phi} = \frac{d\phi}{dt} = \frac{h}{r^2} = h u^2$

$$\frac{dx}{dt} = \frac{d}{dt} \left( \frac{1}{u} \right) = -\frac{1}{u^2} \frac{du}{dt} = -\frac{1}{u^2} \frac{du}{d\theta} \frac{d\theta}{dt} = -\frac{hu^2}{u^2} \frac{du}{d\theta}$$

$$\frac{dr}{dt} = -h \frac{dr}{d\phi}$$



$$\therefore \frac{d^2 x}{dt^2} = \frac{d}{dt} \left( -h \frac{du}{d\theta} \right) = -h \frac{d^2 u}{d\theta^2} \frac{d\theta}{dt}$$

$$\left( \frac{d^2 x}{dt^2} = -h^2 u^2 \frac{d^2 u}{d\theta^2} \right) \begin{cases} h = r^2 \dot{\theta} = \frac{1}{u^2} \\ \Rightarrow \dot{\theta} = hu^2 \end{cases}$$

Put (3) in (1)

$$\frac{1}{u^3} \left[ -h^2 u^2 \frac{d^2 u}{d\theta^2} \right] - h^2 = -\frac{GM}{u}$$

$$\Rightarrow -\frac{h^2}{u} \frac{d^2 u}{d\theta^2} - h^2 = -\frac{GM}{u}$$

$$\Rightarrow \frac{d^2 u}{d\theta^2} + u = \frac{GM}{h^2}$$

$$\Rightarrow \frac{d^2}{d\theta^2} \left[ u - \frac{GM}{h^2} \right] + \left[ u - \frac{GM}{h^2} \right] = 0$$

$$\Rightarrow \boxed{\frac{d^2 u}{d\theta^2} + u = \frac{P}{h^2 u^2}} \quad \text{--- (4)} \quad \begin{cases} P = \frac{\text{force}}{\text{mass}} = \frac{GM}{r^2} \\ F = -\frac{GMm}{r^2} \end{cases}$$

This is the ~~general~~ differential Eqn of the orbit of a particle moving under an attractive central force P per unit mass

$$\boxed{\frac{d^2 u}{d\theta^2} + u = \frac{k}{h^2}} \quad \text{--- (5)} \quad \begin{cases} \forall \text{ let } k = P r^2 = \frac{P}{u^2} \\ k = +GM \end{cases}$$

General Sol<sup>n</sup> of this Eqn is

$$u = \frac{k}{h^2} + A \cos(\theta - \theta_0) \quad \text{--- (6)} \quad \begin{cases} A \text{ \& } \theta_0 \text{ are constants of Eqn.} \end{cases}$$

Also

Let T, U & E be the kinetic, pot. & total energy per unit mass.



9.00

$$\therefore T = \frac{1}{2} v^2 = \frac{1}{2} \left[ \left( \frac{dr}{dt} \right)^2 + r^2 \left( \frac{d\theta}{dt} \right)^2 \right]$$

10.00

$$T = \frac{1}{2} h^2 \left\{ \left( \frac{du}{d\theta} \right)^2 + u^2 \right\}$$

11.00

$$\Rightarrow E - \underset{\substack{\uparrow \\ \text{P.E.}}}{U} = \frac{1}{2} h^2 \left\{ \left( \frac{du}{d\theta} \right)^2 + u^2 \right\}$$

12.00

$$\Rightarrow \left( \frac{du}{d\theta} \right)^2 + u^2 = \frac{2(E-U)}{h^2} \quad \text{--- (7)}$$

1.00

2.00

P.E per unit mass,  $U = -\frac{K}{r} = -K u$  --- (8)

3.00

Put the value of (8) & (6) in eqn. no. (7)

4.00

$$\left( -A \sin(\theta - \theta_0) \right)^2 + \left( \frac{K}{h^2} + A \cos(\theta - \theta_0) \right)^2 = \frac{2}{h^2} \left[ E + K \left( \frac{K}{h^2} + A \cos(\theta - \theta_0) \right) \right]$$

5.00

$$\Rightarrow \left[ A^2 \sin^2(\theta - \theta_0) + \frac{K^2}{h^4} + A^2 \cos^2(\theta - \theta_0) \right] = \left[ \frac{2E}{h^2} + \frac{2K^2}{h^4} + \frac{2KA \cos(\theta - \theta_0)}{h^2} \right]$$

6.00

$$+ 2AK \frac{A \cos(\theta - \theta_0)}{h^2}$$

7.00

$$\Rightarrow A^2 = \frac{K^2}{h^4} + \frac{2E}{h^2} \quad \text{--- (9)}$$

8.00

Putting (9) in (6) ---

9.00

$$u = \frac{K}{h^2} + \sqrt{\frac{K^2}{h^4} + \frac{2E}{h^2}} \cos(\theta - \theta_0) \quad \text{let } \theta_0 = 0$$

$$u = \frac{K}{h^2} \left( 1 + \sqrt{1 + \frac{2Eh^2}{K^2}} \cos \theta \right) \quad \text{--- (10)}$$

Eqn. of Conic Section in polar coordinate is given by

$$u = \frac{1}{\ell} (1 + e \cos \theta) \quad \ell \rightarrow \text{semilatus Rectum}$$

ie. half of focal chord parallel to directrix.

$e$  = eccentricity

Comparing eqn. (10) & (11)

$$l = \frac{h^2}{K}$$

f

$\Rightarrow$

$e =$

$$\sqrt{1 + \frac{2Eh^2}{K^2}}$$

11.00

12.00

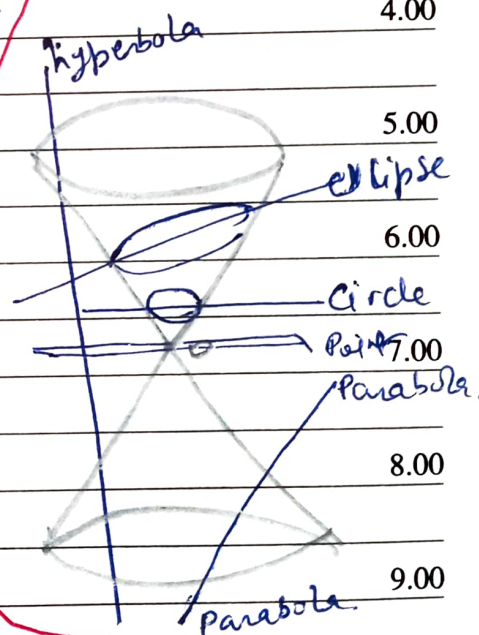
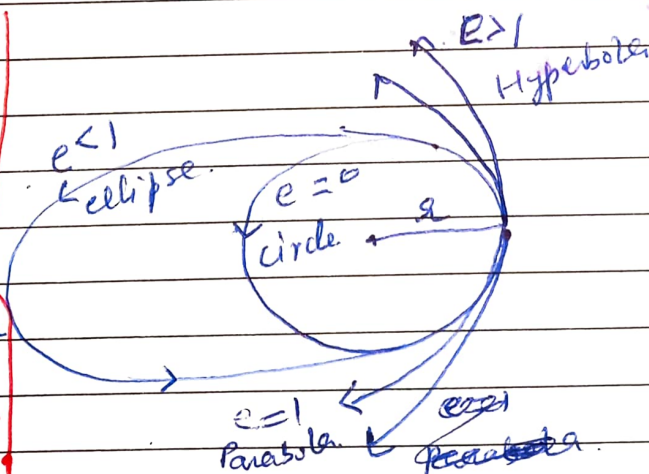
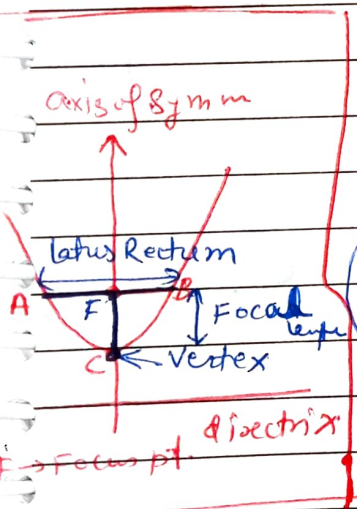
The conic may be of following type:

Conic Section

= Intersection of right circular cones by a plane.

$e$  = measure of noncircularity of an elliptical orbit.

- ① ellipse if  $e < 1$  or  $E < 0$  or -ve.
- ② parabola if  $e = 1$  or  $E = 0$
- ③ hyperbola if  $e > 1$  or  $E > 0$
- ④ circle if  $e = 0$ , or  $E = \frac{-K^2}{2h^2}$



Let  $b \rightarrow$  semi-minor axis of ellipse

$a \rightarrow$  semi-major " " "

such that  $b^2 = a^2(1 - e^2) \Rightarrow e = \sqrt{1 - \frac{b^2}{a^2}}$

$$e = \frac{c}{a}$$

$$e = \sqrt{1 - \left(\frac{b}{a}\right)^2}$$

$$\left\{ \begin{aligned} \therefore l &= \frac{b^2}{a} \\ &= \frac{h^2}{K} \end{aligned} \right.$$

(13)



9.00

Comparing Eqn. (12) + (13)

10.00

$$l = \frac{h^2}{K} \quad \therefore a = \frac{-K}{2E} \Rightarrow$$

11.00

$$h = \sqrt{Ka(1-e^2)}$$

T.E / mass of Particle.

$$E = -\frac{K}{2a}$$

12.00

Speed of the Particle :  $\frac{KE}{mass}$

1.00

$$\frac{\text{Energy}}{\text{man}} = E = \frac{1}{2} v^2 - \frac{K}{2}$$

2.00

$$\rightarrow \frac{-k}{2a} = \frac{1}{2}v^2 - \frac{k}{r}$$

3.00

$$\Rightarrow v^2 = K \left( \frac{2}{r} - \frac{1}{a} \right)$$

4.00

Periodic Time of the particle :  $\frac{ds}{dE} =$  answer.

$$\frac{ds}{dt} = \frac{1}{2} \dot{\theta}^2 = \frac{h}{2}$$

5.00

6.00

∴ Time taken by the ~~body~~ particle to ~~move~~ sweep the entire area of the ellipse.

7.00

$$T = \frac{\text{Total ~~diameter~~ area of the ellipse}}{\text{Areal vel.}}$$

8.00

$$\tau = \frac{\pi a b}{\frac{h}{g}} = \frac{2 \pi a b}{h}$$

20 Sunday

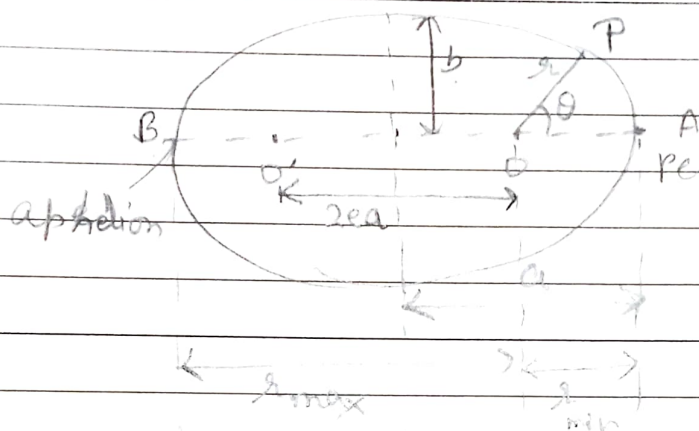
$$\Rightarrow \tau^2 = \frac{4\pi a^2 b^2}{h^2} = \frac{4\pi^2 a^2 \tilde{a}^2 (1-e^2)}{ka(1-e^2)}$$

$$\begin{cases} b^2 = a^2(1-e^2) \\ a(1-e^2) = \frac{h^2}{k_0} \end{cases}$$

$$T^2 = \frac{4\pi^2 a^3}{K} \Rightarrow T = 2\pi \sqrt{\frac{a^3}{K}}$$

$$\Rightarrow \boxed{\tau^2 \propto a^3} \quad (15)$$

$\Rightarrow$  Two planets having same semimajor axis have the same time period.



$e = \frac{\text{distance b/w focus}}{\text{length of major axis}}$

$$\text{Perihelion } e = \frac{OO'}{2a}$$

$$\Rightarrow \boxed{OO' = 2ea}$$

In general

$$\frac{r}{a} = 1 + e \cos \theta$$

$$\left. \begin{aligned} l &= \frac{b^2}{a} \end{aligned} \right\}$$

$\therefore$  at pt A  $\theta = 0$

$$\boxed{\frac{r}{r_{\min}} = 1 + e} \quad (16)$$

$$= \frac{h^2}{k}$$

$$\Rightarrow r_{\min} = \frac{r}{1+e}$$

$$\Rightarrow \boxed{r_{\min} = \frac{b^2/a}{1+e}}$$

at pt B,  $\theta = 180^\circ$

$$\boxed{\frac{r}{r_{\max}} = 1 - e} \quad (17)$$

$$\Rightarrow r_{\max} = \frac{r}{1-e}$$

$$\boxed{r_{\max} = \frac{b^2/a}{1-e}}$$

$$\text{eccentricity } e = \frac{r_{\max} - r_{\min}}{r_{\max} + r_{\min}}$$



$$a = \frac{r_{\max} + r_{\min}}{2} \quad \left\{ \begin{array}{l} b^2 = a^2(1 - e^2) \\ l = \frac{b^2}{a} \end{array} \right.$$

a avg of max & min

distance of the planet from centre of force  
ie the sun.

NOTE

At the perihelion & Aphelion, the vel. vector is  $\perp$  to the position vector  $\vec{r}$ .

At these points

$$v = r\dot{\theta}$$

$$= r h \frac{1}{r^2}$$

$$= \frac{h}{r}$$

$$v = \frac{J}{mr}$$

$$\left\{ \begin{array}{l} v_1 = 0 \\ v_2 = r\dot{\theta} \end{array} \right.$$

$$\left\{ \begin{array}{l} h = r^2 \dot{\theta} \end{array} \right.$$

$$\left\{ \begin{array}{l} h = \frac{J}{m} \end{array} \right.$$

$$\left\{ \begin{array}{l} \frac{ds}{dt} = \frac{1}{2} r^2 \dot{\theta} \end{array} \right.$$

$$\left\{ \begin{array}{l} = \frac{1}{2} r v \end{array} \right.$$

$$\left\{ \begin{array}{l} = \frac{1}{2} \frac{J}{m} \end{array} \right.$$

∴ Total Energy of the system

$$E = \frac{1}{2} m \left( \frac{J}{mr} \right)^2 - \frac{GMm}{r}$$

$$= \frac{1}{2} \frac{J^2}{m} \left( \frac{1}{r_{\min}} \right)^2 - GMm \left( \frac{1}{r_{\min}} \right) \quad \text{--- (18)}$$

$$= \frac{1}{2} \frac{J^2}{m} \left( \frac{1}{r_{\max}} \right)^2 - GMm \left( \frac{1}{r_{\max}} \right) \quad \text{--- (19)}$$

~~Let the value of  $r_{\min}$  from (16) to (18)~~

$$\text{e Now } l = \frac{h^2}{GM} = \frac{J^2}{GMm^2} \quad \left\{ \begin{array}{l} \therefore h = \frac{J}{m} \end{array} \right.$$

$$\frac{1}{r_{min}} = \frac{(1+e) G M m^2}{J^2}$$

$$E = \frac{1}{2} \frac{J^2}{m} \left[ \frac{(1+e) G M m^2}{J^2} \right]^2 - G M m \left[ \frac{(1+e) G M m^2}{J^2} \right]$$

$$= \frac{1}{2} \frac{G^2 M^2 m^3}{J^2} (1+e)^2 - \frac{G^2 M^2 m^3}{J^2} (1+e)$$

$$= \frac{G^2 M^2 m^3}{J^2} (1+e) \left[ \frac{1+e}{2} - 1 \right]$$

$$= \frac{G^2 M^2 m^3}{J^2} (1+e) \left( \frac{e-1}{2} \right)$$

$$E = - \frac{G^2 M^2 m^3}{2 J^2} (1-e^2) \quad \text{for } e=0 \quad E = - \frac{G^2 M^2 m^3}{2 J^2}$$

$\Rightarrow$  if  $e < 1$  The path of planet is elliptical.

NOTE: from selection derived earlier

$$e = \sqrt{1 + \frac{2 E L^2}{K^2}}$$

$$K = -G M$$

$$h = \frac{J}{m}$$

$$\Rightarrow \frac{E}{m} = \frac{(e^2 - 1) K^2}{2 h^2}$$

$$= \frac{(e^2 - 1) G^2 M^2 m^2}{2 J^2}$$

$$\frac{E}{m} = - \frac{(1-e^2) G^2 M^2 m^2}{2 J^2}$$

$$\Rightarrow E = - \frac{(1-e^2) G^2 M^2 m^3}{2 J^2}$$

in General  $E = - \frac{(1-e^2) G^2 M^2 m^2}{2 J^2} h$

$\mu$  = reduced mass of the system